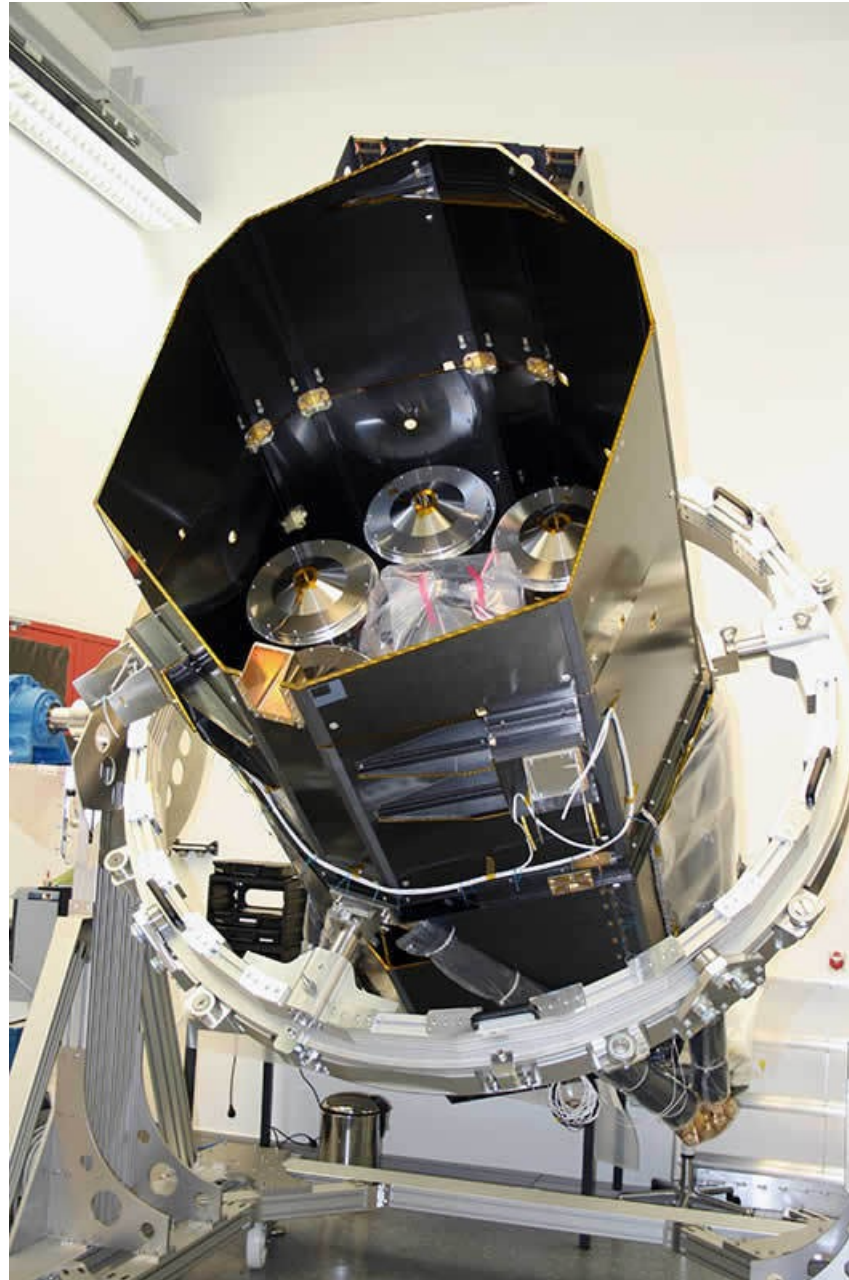




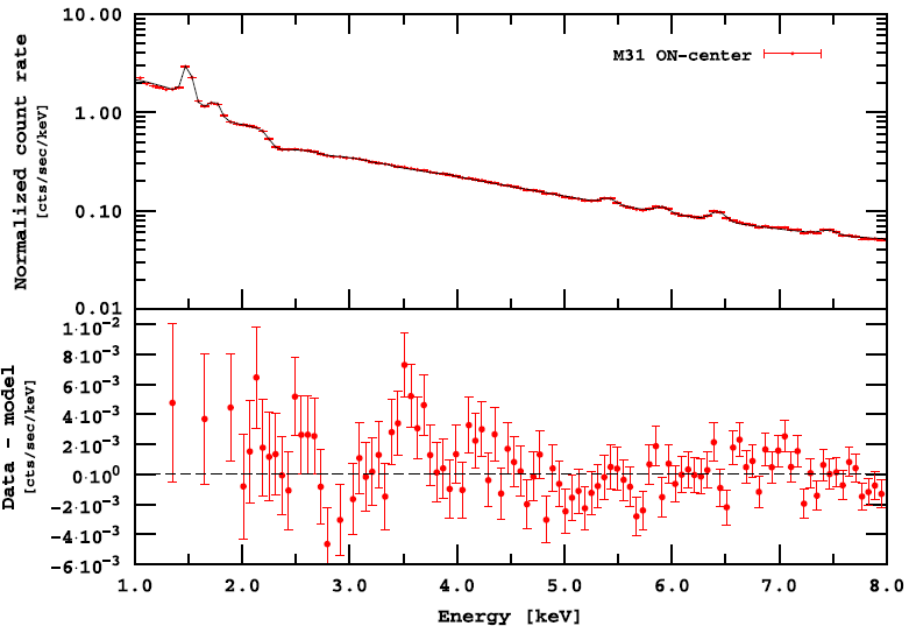
Decaying dark matter search with eROSITA



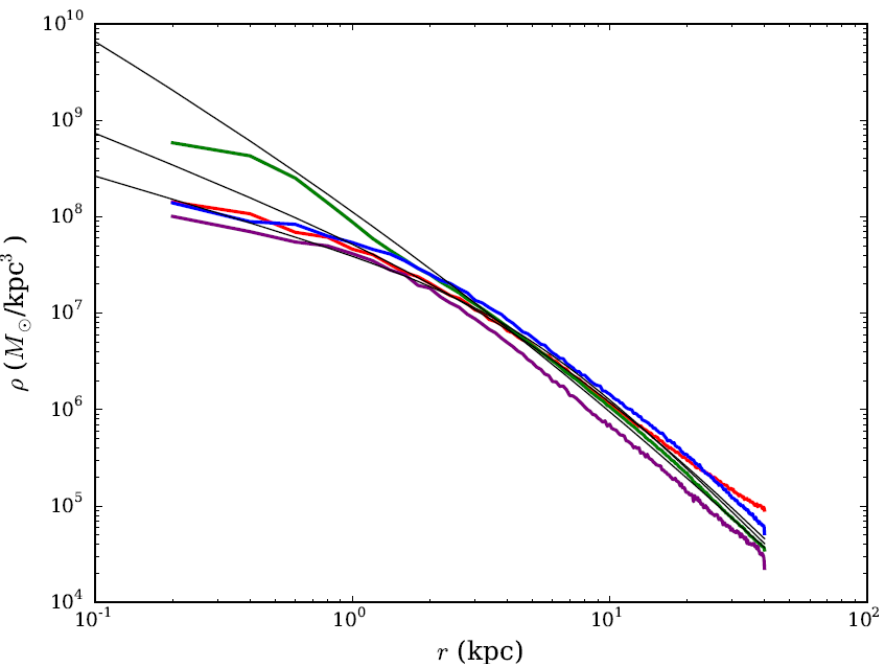
D. Malyshev
(IAAT, Tübingen University)



General remarks...



3.55 keV line in M31 from Boyarsky et al (1402.4119)



Uncertainties in DM density in LMC, from Buckley et al 1502.01020

– In astrophysics there is ***no way to directly measure*** the parameters (e.g. mass) of DM particle.

– The only option is to ***deduce*** these parameters ***from*** the ***observations*** that do not have convenient astrophysical explanations.

– The most straightforward way to detect decaying/annihilating dark matter is to ***observe*** the corresponding feature in DM-dominated objects' ***(photon) spectra***. The signal depends on the amount (and spatial distribution in annihilating DM case) of DM in instrument's FoV

Caveats:

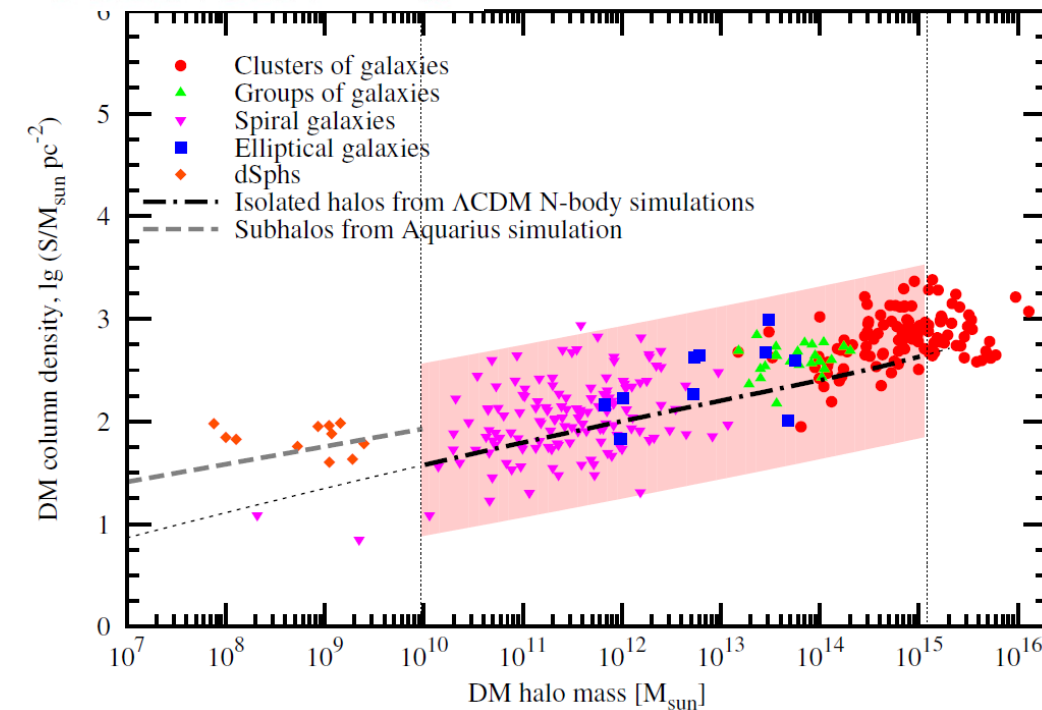
– DM profiles can have ***large uncertainties***, especially near the centers of DM dominated object, from where the most of the DM signal is naturally expected.

– DM decay/annihilating spectral feature can be confused with the ***convenient astrophysical line/feature***. Or with the feature in instrument's effective area.

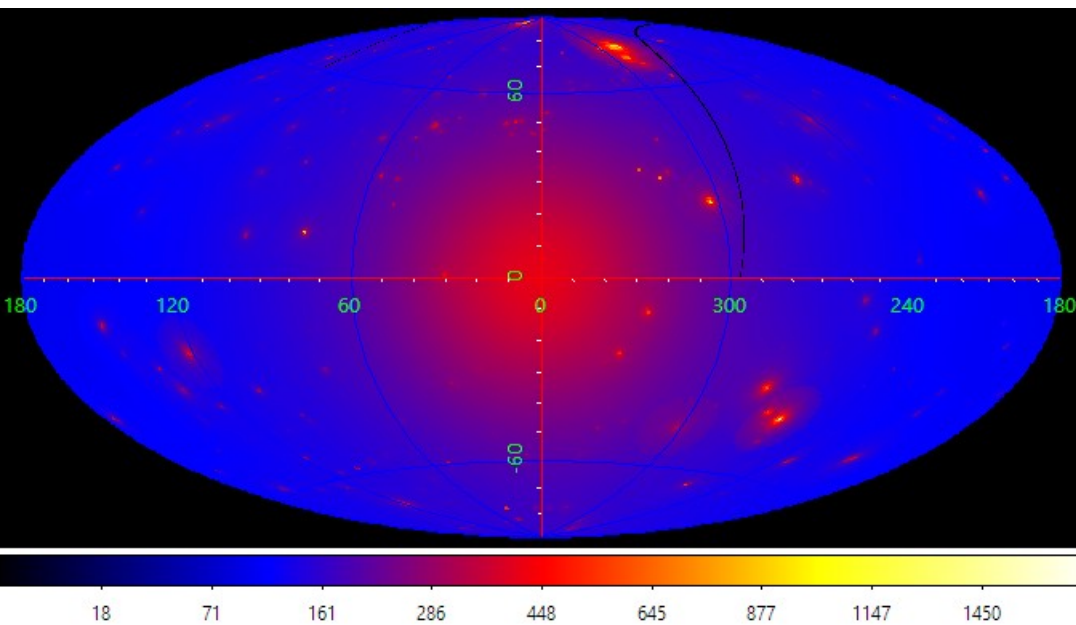
– The expected feature is generally weak and requires the analysis of a large amount of the data. At this step one can encounter previously ***unknown instrument systematic*** effects.



General remarks...



From Boyarsky et al 0911.1774



Where to search for decaying dark matter?

- Our galaxy?
- Clusters?
- Dwarf spheroidals?
- something else?

The expected flux F from DM decay (annihilation) is proportional to the DM column density

$$dF \sim S d\Omega \quad S_{\text{dec}} \sim \int \rho d\ell$$

Note $d\Omega$ in the formula for the expected signal. If FoV of the instrument is much larger than the size of the object – the contribution of the object to the signal will not be significant!

If the size of the object is comparable with the FoV – it appears that **S is compatible for different types of objects**, from dSphs to Clusters. Other considerations (e.g. lower astrophysical background ; quality of the knowledge of DM density distribution) should be invoked in this case.



Sterile neutrino Dark Matter

SM

mass →	2.4 MeV	1.27 GeV	171.2 GeV
charge →	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$
name →	u	c	t
	Left up Right	Left charm Right	Left top Right
Quarks	4.8 MeV	104 MeV	4.2 GeV
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$
	d	s	b
	Left down Right	Left strange Right	Left bottom Right
Leptons	0 eV	0 eV	0 eV
	0	0	0
	ν_e	ν_μ	ν_τ
	Left electron Right	Left muon Right	Left tau Right
	0.511 MeV	105.7 MeV	1.777 GeV
	-1	-1	-1
	e	μ	τ
	Left electron Right	Left muon Right	Left tau Right

nuMSM

mass →	2.4 MeV	1.27 GeV	171.2 GeV
charge →	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$
name →	u	c	t
	Left up Right	Left charm Right	Left top Right
Quarks	4.8 MeV	104 MeV	4.2 GeV
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$
	d	s	b
	Left down Right	Left strange Right	Left bottom Right
Leptons	<0.0001 eV	~0.01 eV	~0.04 eV
	0	0	0
	ν_e	ν_μ	ν_τ
	Left electron Right	Left muon Right	Left tau Right
	0.511 MeV	105.7 MeV	1.777 GeV
	-1	-1	-1
	e	μ	τ
	Left electron Right	Left muon Right	Left tau Right

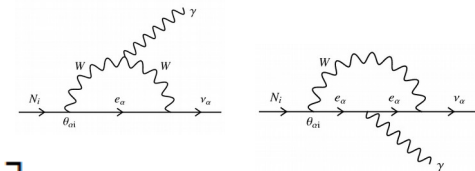
– ν MSM Dark Matter

All particles in the SM except neutrino can have positive («right handed») and negative («left-handed») projection of the spin onto their momentum. Standard neutrinos can be only left-handed. ν MSM proposes to add right-handed «sterile» companions to the neutrinos.

NuMSM is nice because:

- it is a «natural» extension of the SM
- It can explain low neutrino masses (seesaw mechanism) / neutrino oscillations
- it provides a way to explain baryon number generation in the Early Universe
- gives candidate(s) for the Dark Matter
- probably accounts for some observed phenomena (e.g. pulsar kicks)

There are straightforward ways to observe decays of the sterile neutrino.
Radiative decay channel suggests the narrow line at energy $E = M_s/2$



$$F_{DM} \approx 1.4 \times 10^{-7} \left[\frac{\sin^2(2\theta)}{10^{-11}} \right] \left[\frac{m_{DM}}{10 \text{ keV}} \right]^4 \left[\frac{d}{100 \text{ kpc}} \right]^{-2} \left[\frac{M_{DM, FoV}}{10^7 M_\odot} \right] \frac{\text{ph}}{\text{cm}^2 \text{ s}}$$

Note: the signal for decaying DM does not depend on the exact DM distribution profile, but just on a total mass within instrument's FoV!





Sterile neutrino Dark Matter

Non-spectral bounds

– **Tremaine-Gunn bound** (phase space density limit)

Since sterile neutrinos are fermions it is not possible to put too much particles within some DM-dominated objects (e.g. dSphs) within allowed velocity dispersion range.

See S. Tremaine & J. E. Gunn, Phys. Rev. Lett. 42 (1979) 407.

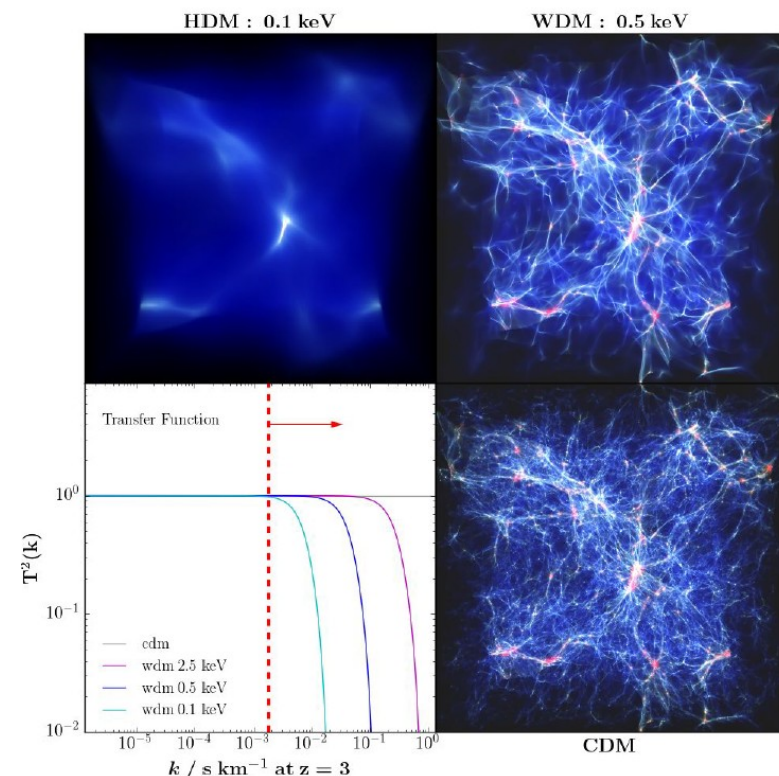
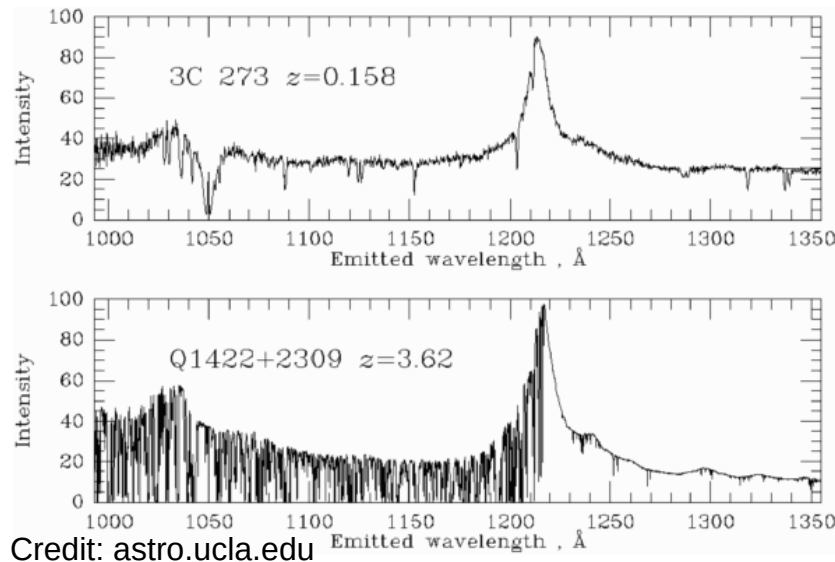
Implies limit $m \geq 1 \text{ keV}$

– **Lyman-alpha bound**

High-z AGN light travelling through the Universe can be absorbed and re-radiated by the neutral hydrogen all way to us. The frequency of re-radiated photon is const in neutral hydrogen frame and z-shifted for us. Absorption and re-radiation lead to the numerous (z-shifted) Ly-alpha lines.

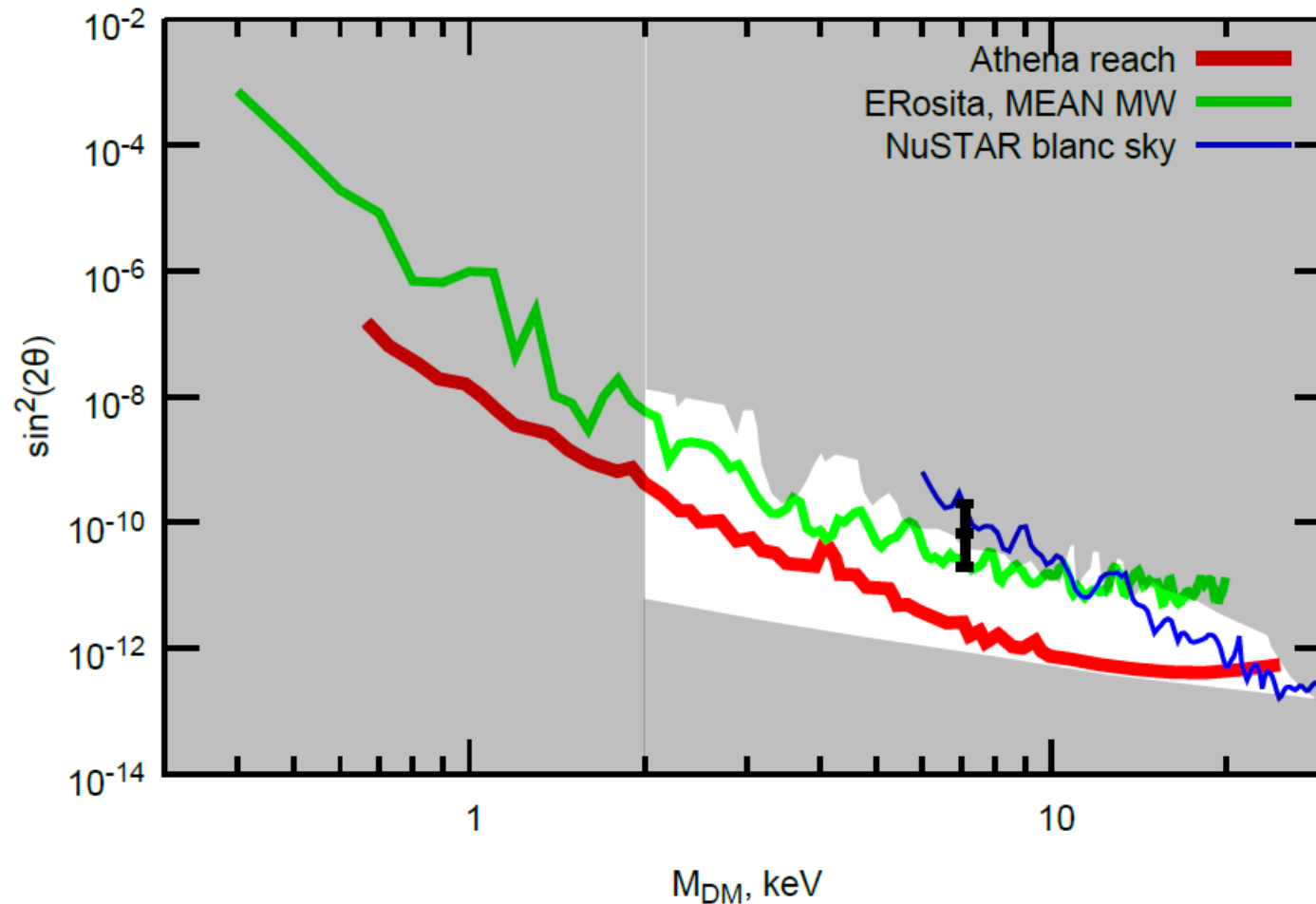
Neutral hydrogen is a tracer of DM distribution. Power spectrum of the DM distribution is a function of DM particle mass and can be compared with the deduced from neutral hydrogen distribution.

Implies limit $m \geq 2..4 \text{ keV}$





Sterile neutrino Dark Matter

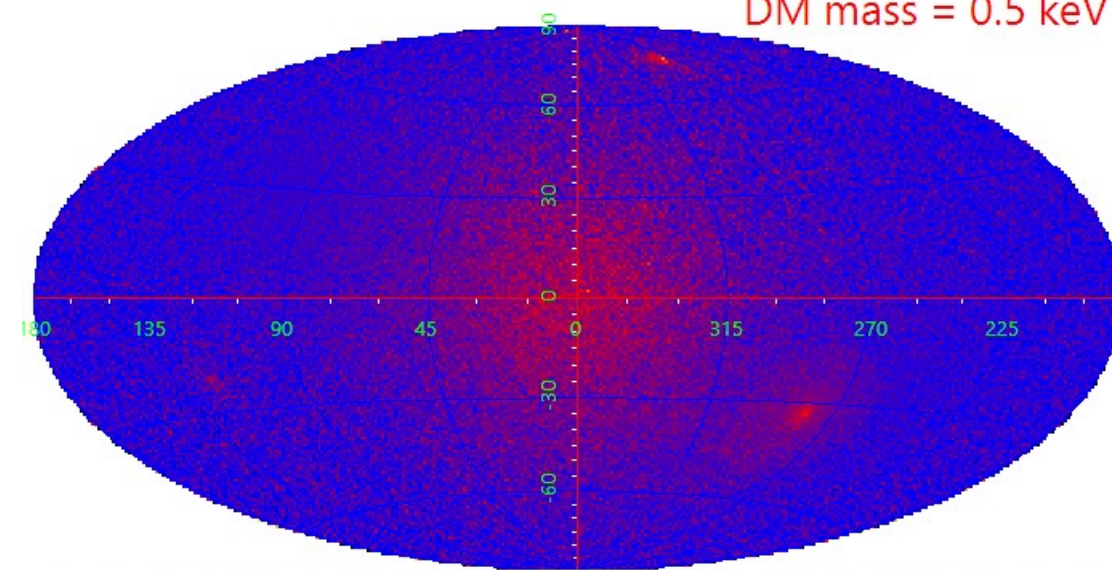


eROSITA limits on decaying sterile neutrino dark matter. The estimations for 1Msec observations of Segue I dSph (similar to Athena reach limits).



Sterile neutrino Dark Matter

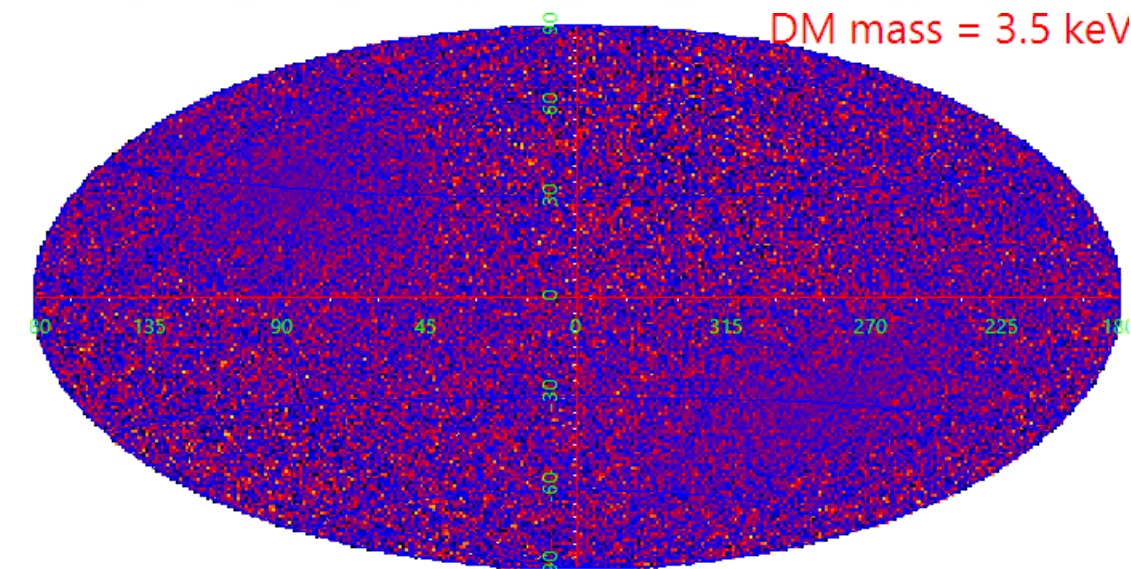
DM mass = 0.5 keV



eROSITA view of decaying sterile neutrino dark matter for $m_{\text{DM}} = 0.5 \text{ keV}$ and $m_{\text{DM}} = 3.5 \text{ keV}$ (assuming the DM parameters from existing upper bounds)

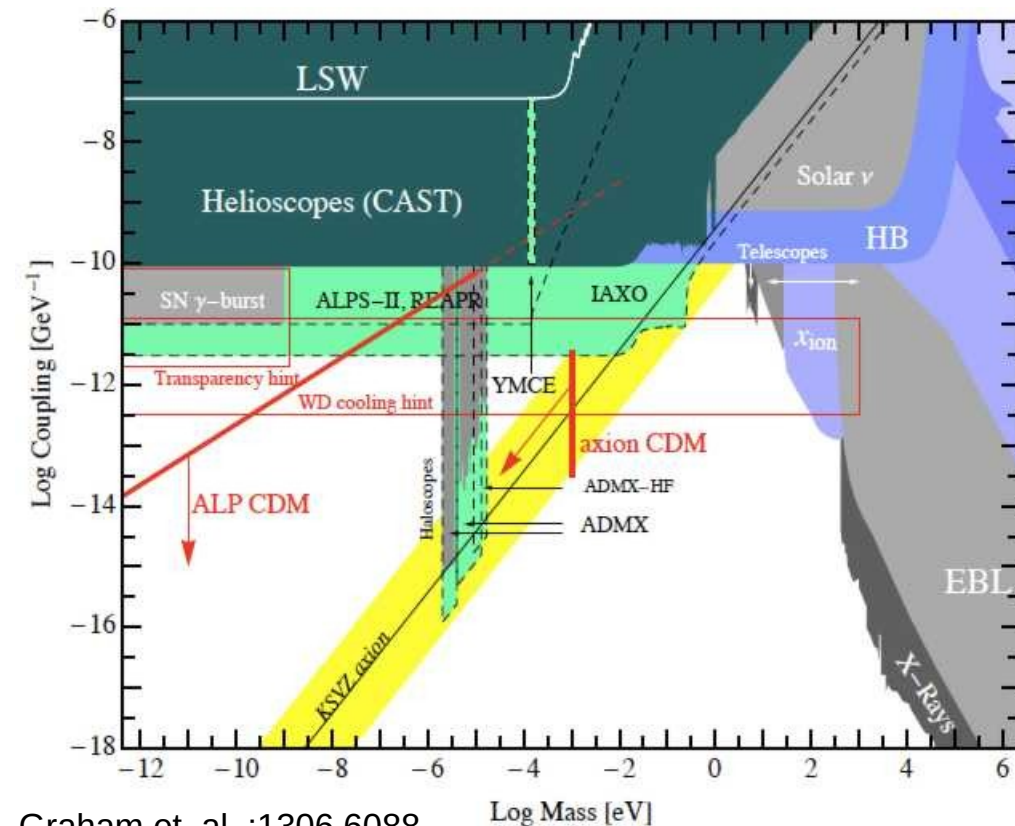
DM is clearly seen for energies below $\sim 1.5 \text{ keV}$. Possibility for ON-OFF analysis and/or for the correlation studies.

DM mass = 3.5 keV





Axion Dark Matter



Graham et. al. :1306.6088

$$\Gamma_{a \rightarrow \gamma\gamma} = \frac{g_{a\gamma\gamma}^2 m_a^3}{64\pi} = 7.6 \cdot 10^{-26} \left(\frac{g_{a \rightarrow \gamma\gamma}}{10^{-10} \text{GeV}^{-1}} \right)^2 \left(\frac{m_a}{1 \text{eV}} \right)^3$$

$$P_{\gamma \rightarrow a} = (\Delta_{a\gamma s})^2 \frac{\sin^2(\Delta_{osc s}/2)}{(\Delta_{osc s}/2)^2}$$

$$\Delta_{osc}^2 = (\Delta_{pl} - \Delta_a)^2 + 4\Delta_{a\gamma}^2$$

$$\Delta_{a\gamma}/\text{Mpc}^{-1} = 0.15 \cdot \frac{g_{a\gamma\gamma}}{10^{-10} \text{GeV}^{-1}} B_{nG}$$

$$\Delta_a/\text{Mpc}^{-1} = -7.7 \times 10^{28} \left(\frac{m_a}{1 \text{eV}} \right)^2 \left(\frac{\omega}{1 \text{eV}} \right)^{-1}$$

$$\Delta_{pl}/\text{Mpc}^{-1} = -11.1 \left(\frac{n_e}{10^{-7} \text{cm}^{-3}} \right) \cdot \left(\frac{\omega}{1 \text{eV}} \right)^{-1}$$

Mirizzi et.al.:astro-ph/0607415

– Axion

Considered to solve strong CP problem in QCD and a good candidate to be DM, see e.g. Marsh, 1510.07633 for the review.

Despite the low mass of the axion, there are ways to make it cold – «misalignment mechanism»

Can be detected via two photon decay and conversion to photon (and v.v.) in a magnetic field (Primakoff effect).

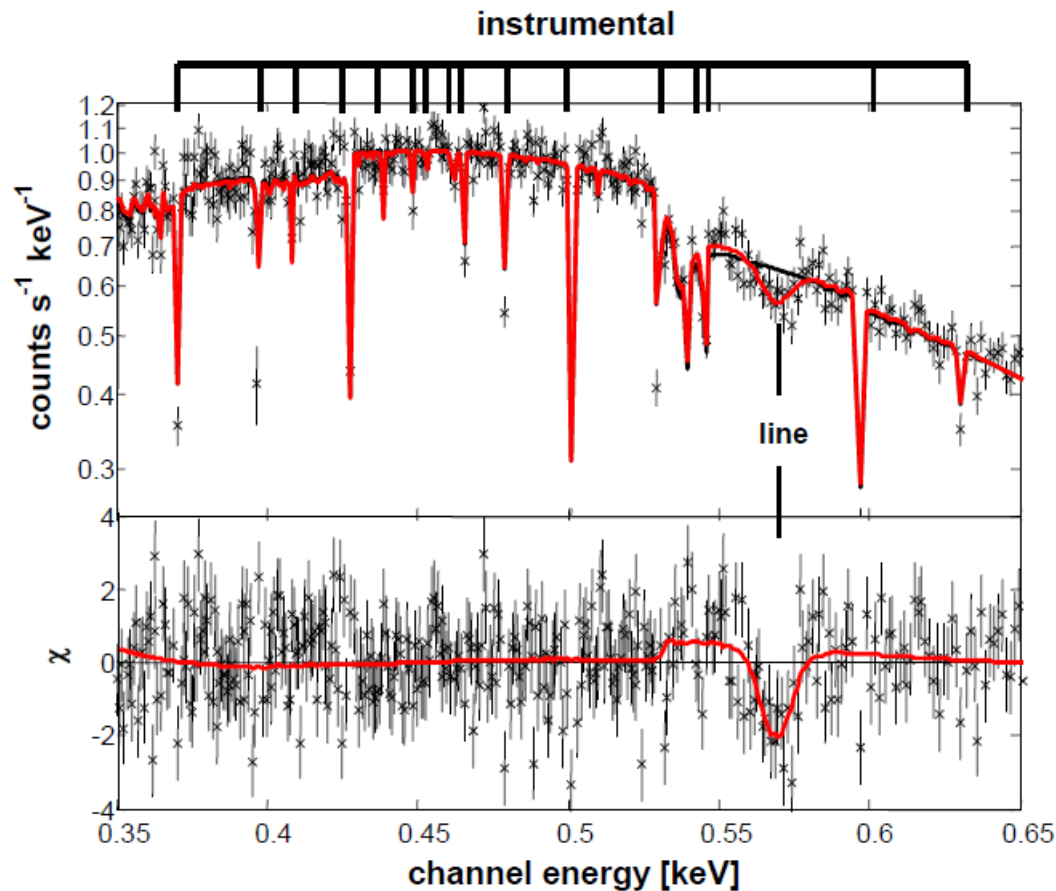
Astrophysical consequences:

- additional energy loss channel for stellar/SN evolution
- Features in pulsars/magnetars spectra
- Distant sources dimming
- Additional polarization of distant sources' light.

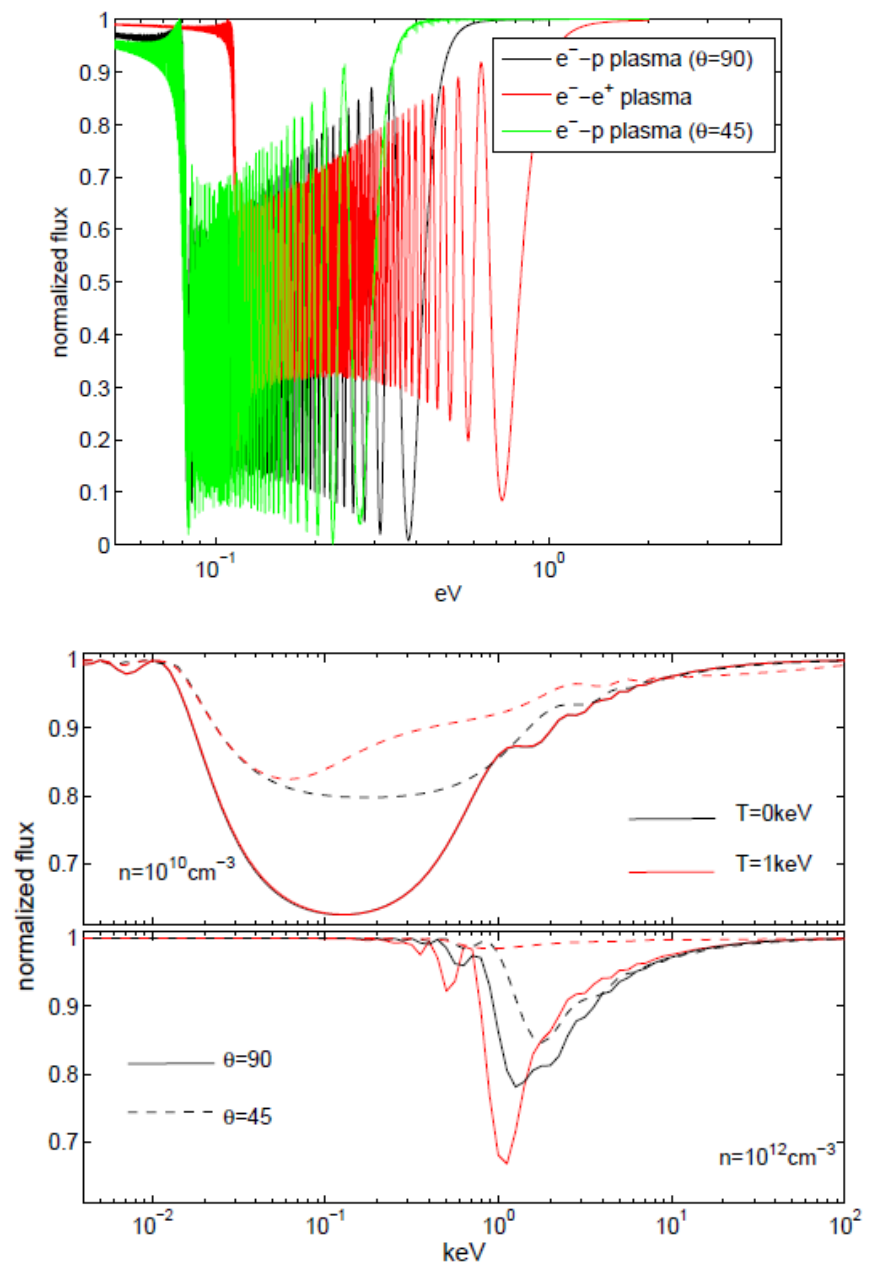
Opens a window for direct search for axions: ADMX, CAST, ALPS.



Axion Dark Matter



XMM/RGS spectrum of RXJ0720.4-3125
 (adopted from 1109.2506)



Features in spectra of pulsars (top) and X-ray binaries from axion-photon conversion. Adopted from 0806.0411



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